Impact of Radiation on MHD Oscillatory Convective Flow of Heat Absorbing Visco - Elastic Dusty Fluid Confined In Horizontal Channel

M.Shanthi[1], Dr.P.T.Hemamalini[2]

[1] Department of Mathematics, Karpagam College of Engineering, TamilNadu,India,shanthi2saba@gmail.com

[2]Department of Mathematics, Karpagam Academy of Higher Education, TamilNadu,India,hema_0869@yahoo.co.in

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Abstract: The unsteady freeconvective magnetohydrodynamic flow of an conducting visco elastic dusty fluid with transverse magnetic field and radiation confined by horizontal channel is considered. The influence of oscillatory pressuregradient and the movement of the top plate are incorporated. The heat generated is established in both motions of fluid and dust particles and is high enough to radiate heat. The solutions of velocity distribution, temperature distribution of the fluid, velocity and energy of the dust particle are achieved by employing perturbation scheme. The physical parameters involved in both profiles for fluid and dust particles are examined with aid of graph.

Keywords: MHD, Horizontal channel, radiation , visco elastic dusty fluid.

1. Introduction

In the view of several physical problems the research of porous media in magnetic field and radiation plays major role. extensive research has been carried out and it has paid the way to study the magnetohydrodynamic flow in channel together with the magnetic field and radiation .

Ahmed, et al, studied about the thermal diffusion effect on a three-dimensional MHD free convection with mass transfer flow from a porous vertical plate. Venkatesh and Kumara, discussed about the exact solution of an unsteady conducting dusty fluid flow between non-tortional oscillating plate and a long wavy wall. Closed form solutions for unsteady free convection flow of a second grade fluid over an oscillating vertical plate is explained by Ali, et al. Dey learnt about dusty hydromagnetic oldryod fluid flow in a horizontal channel with volume fraction and energy dissipation. Sheikh, et al, investigated by comparison and analysis of the Atangana – Baleanu and Caputo – Fabrizio fractional derivatives for generalized Casson fluid model with heat generation and chemical reaction. Dielectrophoretic choking phenomenon of a deformable particle in a converging-diverging microchannel is analyzed by Zhou, T. et al. Dey, explained viscoelastic fluid flow through an annulus with relaxation, retardation effects and external heat source/sink. Bilal, et al, discussed about two-phase fluctuating flow of dusty viscoelastic fluid between non conducting rigid plates with heat transfer. Zhou, et al. learnt AC dielectrophoretic deformable particle-particle interactions and their relative motions. Khan, et al. discussed about effects of relative magnetic field, chemical reaction, heat generation and Newtonian heating on convection flow of casson fluid over a moving vertical plate embedded in a porous medium. Rajakumar, et al, explained about the influence of dufour and thermal radiation on unsteady mhd walter's liquid model-b flow past an impulsively started infinite vertical plate embedded in a porous medium with chemical reaction, hall and ion slip current.

2. Physical Description of the problem

The unsteady oscillating free convective flow of heat absorbing viscoelastic dusty fluid in horizontal plates has been considered inclusive of transverse magnetic field and radiation effect. The movement of the top plate with free stream velocity U (t) induced by the oscillating pressure gradient.

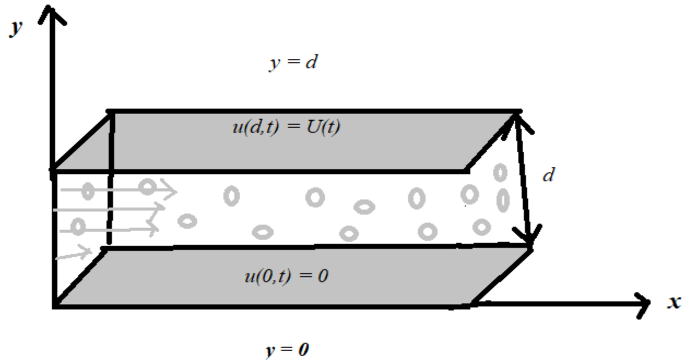

Figure 1 Physical Configuration

The flow is considered along the X- axis. The plate at y = d fluctuating with freestream velocity U (t) while the plate at y = 0 is at rest. The fluid and dust particles velocities are given by u and v.The temperature between the two plates are high enough to radiate heat. The lower plate is marked with ambient temperature T_∞ and other plate temperature is sustained with T_w whereas Tp denotes the temperature of the particle

$$\frac{\partial u}{\partial t}=\frac{1}{\rho}\frac{\partial p}{\partial x}+\left(\upsilon+\frac{\alpha_1}{\rho}\frac{\partial}{\partial t}\right)\frac{\partial^2 u}{\partial y^2}-\frac{K_0N_0}{\rho}(u-v)-\frac{\sigma B_0^2 u}{\rho}+g\beta_T(T-T_\infty) \quad (1)$$

$$m\frac{\partial}{\partial t}v=(u-v)K_0 \quad (2)$$

$$\frac{\partial T}{\partial t}=\frac{k}{\rho c_p}\frac{\partial^2 T}{\partial y^2}+\frac{\rho_p C_s}{\rho C_p \gamma_T}(T_p-T)-\frac{Q_0 T}{\rho C_p}-\frac{1}{\rho}\frac{\partial q_r}{\partial y} \quad (3)$$

$$\frac{\partial T_p}{\partial t}=\frac{1}{\gamma_T}\left(T-T_p\right) \quad (4)$$

The relavent boundary conditions are

$$\begin{aligned} &For\, t\le 0, T(y,0)=T_\infty, u(y,0)=0 \\ &t>0, u(0,t)=0, T(0,t)=T_w \qquad at \qquad y=0 \\ &t>0, u(d,t)=U(t), T(d,t)=T_\infty \qquad at \qquad y=d \end{aligned} \quad (5)$$

Consider $v(y,t)=\frac{v_0(y)}{e^{-i\omega t}}$,in order to calculate the velocity of the dust particle and by substituting the value of v in equation (2),

$$v(y,t)=(\frac{K_0}{mi\omega+K_0})u(y,t) \quad (6)$$

Incorporating equation (6) in (1) and at free stream area equation (1) will take the form

$$\frac{\partial u}{\partial t}=-\frac{1}{\rho}\frac{\partial p}{\partial x}+\left(\upsilon+\frac{\alpha_1}{\rho}\frac{\partial}{\partial t}\right)\frac{\partial^2 u}{\partial y^2}-+\frac{K_0N_0(U-u)}{\rho}(\frac{K_0}{mi\omega+K_0}-1)-\frac{\sigma B_0^2}{\rho}(u-U)+g\beta_T(T-T_\infty) \quad (7)$$

For non-dimensionalizating the above equations,the following are considered

$$u^* = \frac{U}{U_0}, y^* = (yd^{-1}),\ t^* = \frac{U_0}{d}t, \theta = -\frac{[T_\infty - T]}{T_w - T_\infty}, \theta_p = \frac{-T_p + T_\infty}{-T_w + T_\infty} \tau^* = \frac{1}{\mu\upsilon}\tau d^2$$

$$\mathrm{Re} = \frac{u_0 d}{\upsilon}, K_1 = \frac{K_0 N_0 d^2}{\rho\upsilon}, \alpha = \frac{\alpha_1 u_0}{\rho\upsilon d}, K_2 = \frac{K_0^2 N_0^2}{\rho\upsilon(mi\omega + K_0)}, M = \frac{\sigma B_0^2 d^2}{\rho\upsilon},$$

$$Gr = \frac{g\beta_T (T_w - T_\infty)}{\upsilon u_0}, \phi = \frac{dQ_0}{\rho C_p}$$

(8)

The equation of momentum, energy of the fluid and dust particles in dimensionless form is given as:

$$R_e \frac{\partial}{\partial t} u = R_e \frac{dU}{dt} + \frac{\partial^2 u}{\partial y^2} + \alpha \frac{\partial}{\partial t}\left(\frac{\partial^2 u}{\partial y^2}\right) + (K_2 - K_1)(U - u) - M(u - U) + G_r \theta$$

(9)

$$\frac{\partial \theta}{\partial t} = (Pe)^{-1} \frac{\partial^2}{\partial y^2}\theta + R\frac{1}{Pe}(\theta_p - \theta) - (\phi + N)\theta \tag{10}$$

$$\frac{\partial \theta_p}{\partial t} = \gamma(\theta - \theta_p) \tag{11}$$

Let assume that $\theta_{(p)} = \left[\frac{\gamma}{i\omega + \gamma}\right]\theta$, then energy equation becomes

$$\frac{\partial \theta}{\partial t} = (pe)^{-1} \frac{\partial^2 \theta}{\partial y^2} + R\frac{1}{Pe}\left[\frac{1}{i\omega + \gamma}\gamma - \theta\right] - (\phi + N)\theta$$

(12)

The corresponding boundary conditions are

$$u(1,t) = U(t), \theta(0,t) = 1, u(0,t) = 0, \theta(1,t) = 0,$$

$$U(t) = 1 + \frac{\in}{2}\left(e^{i\omega t} + e^{-i\omega t}\right) \tag{13}$$

Considering the periodic solution in order to solve the equations

$$u(y,t) = u_0(y) + \in \frac{1}{2}(u_1(y) + u_2(y))e^{i\omega t}$$

$$\theta(y,t) = \theta_0(y) + \in \theta_1(y)e^{i\omega t} \tag{14}$$

Adopting the assumpsions and boundary conditions in equation and momenteum,the harmonic and non harmonic parts are given by

$$\theta_1(y) = 0$$

$$\theta_0(y) = \frac{\sinh\left[\sqrt{m_1} - \sqrt{m_1}y\right]}{\sinh(\sqrt{m_0})} \tag{15}$$

$$\frac{\partial u}{\partial t}\mathrm{Re} = \frac{dU}{dt}\mathrm{Re} + \frac{\partial^2}{\partial y^2}u + \alpha\frac{\partial^3 u}{\partial t \partial y^2} - (K_1 - K_2) - M(u - U) + Gr\left[\frac{\sinh\left[\sqrt{m_1} - \sqrt{m_1}y\right]}{\sinh(\sqrt{m_0})}\right] \tag{16}$$

Using equation (13),(14) in (16)

$$u_1(y) = \frac{\sinh \sqrt{m_2}y - \sqrt{m_2}}{\sinh(\sqrt{m_2})} + 1 \tag{17}$$

$$u_2(y) = \frac{\sinh \sqrt{m_3} y - \sqrt{m_3}}{\sinh(\sqrt{m_3})} + 1$$

(18)

$$u_0(y) = -B\frac{\sinh \sqrt{D} - \sqrt{Dy}}{\sinh(\sqrt{D})} + 1 + (B-1)\frac{\sinh \sqrt{m_1} - \sqrt{m_1 y}}{\sinh(\sqrt{m_1})}$$

(19)

$$u(y,t) = -B\frac{\sinh \sqrt{D} - \sqrt{Dy}}{\sinh(\sqrt{D})} + 1 + (B-1)\frac{\sinh \sqrt{m_1} - \sqrt{m_1 y}}{\sinh(\sqrt{m_1})} + \frac{\in}{2}\left[\begin{array}{l}\left(\frac{\sinh \sqrt{m_2} y - \sqrt{m_2}}{\sinh(\sqrt{m_2})} + 1\right)e^{i\omega t} + \\ = \frac{\sinh \sqrt{m_3} y - \sqrt{m_3}}{\sinh(\sqrt{m_3})} + 1\end{array}\right]$$

(20)

where

$$m_1 = \frac{N(\gamma + i\omega) + R_1(\gamma + i\omega) + Pe\gamma\phi - (\gamma + i\omega)\phi}{\gamma + i\omega}, \quad m_1 = \frac{\mathrm{Re}\, i\omega - D}{1+\alpha}, \quad m_2 = \frac{\mathrm{Re}\, i\omega + D}{1-\alpha}$$

$$D = -(K_2 - K_1) - M, \quad B = 1 + Gr(\frac{1}{D - m_1})$$

3. Results:

The predominent physical parameters involved in the flow are discussed using graph.

Fig 2,5,6 and 7 illustrates about the behaviour of velocity graph of fluid for different values of ϕ , M, Pe and R. For the increase in these parameters ,the fluid velocity profile degrades. whereas fig 3 and 4 depicts about the progress of velocity graph of fluid for incrementing values of Gr and K .

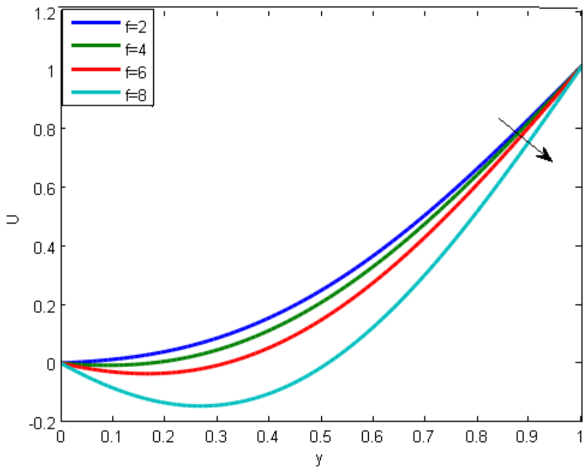

Fig 2: Velocity profile U for ϕ, [Gr=2,M=2,Pe=2,Re=1,K=1]

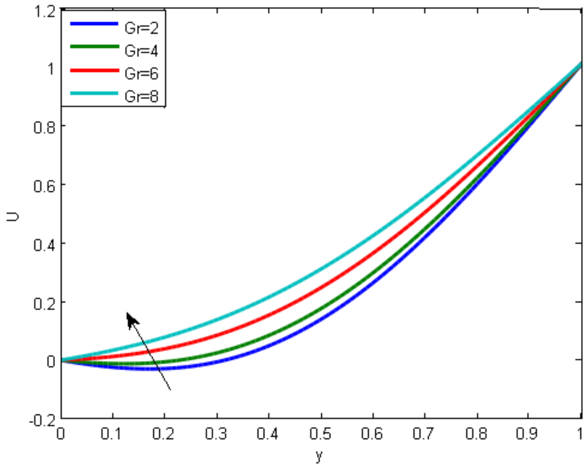

Fig 3: Velocity profile U for Gr, [ϕ =1, M=2,Pe=2,Re=1,K=1]

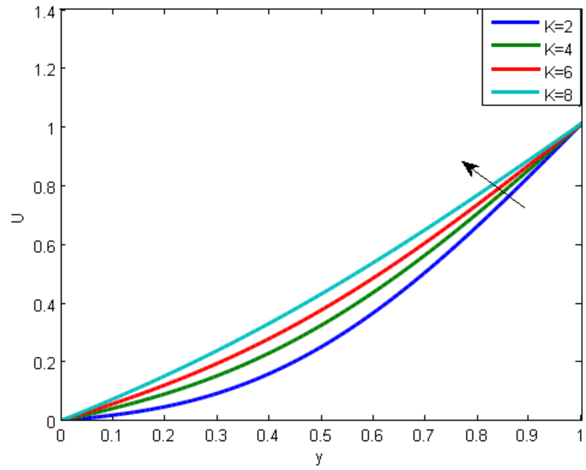

Fig 4: Velocity profile U for K, [ϕ =1, M=2, Pe=2, Re=1, Gr=2]

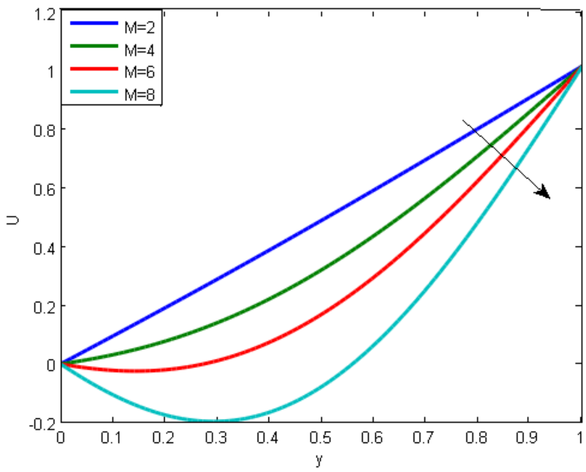

Fig 5: Velocity profile U for M, [ϕ =1,Gr = 2, Pe=2,Re=1,K=1]

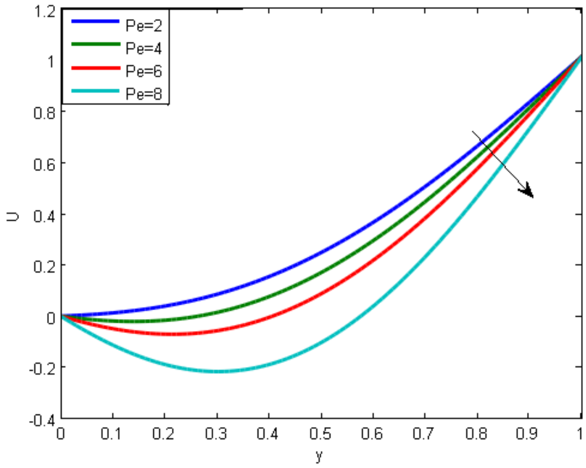

Fig 6: Velocity profile U for Pe, [ϕ =1, M=2, Gr=2,Re=1,K=1]

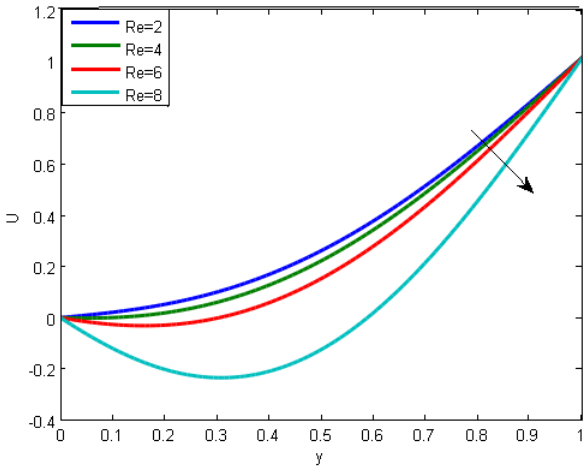

Fig 7: Velocity profile U for Re, [ϕ =1,Gr = 2, M=2, Pe=1,K=1]

Figure 8 - 13 represents the influence of velocity of the dusty particle for various values of R,Pe,M, ϕ ,Gr and m. The velocity profile degrades with increasing these physical factors except for the parameter m.On incrementing the values of m,it leads to the causes the upgrade in the velocity of dust particles.

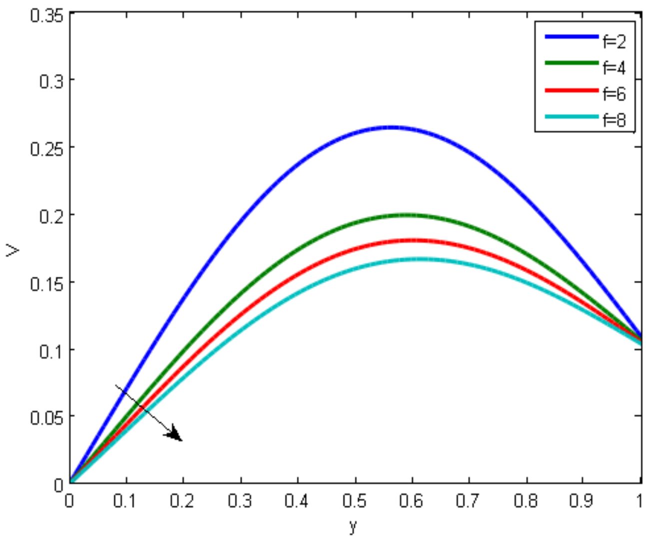

Fig 8: Particle velocity profile V for ϕ, [M=2,Gr = 2, Pe=2,Re=1,K=1,m=1]

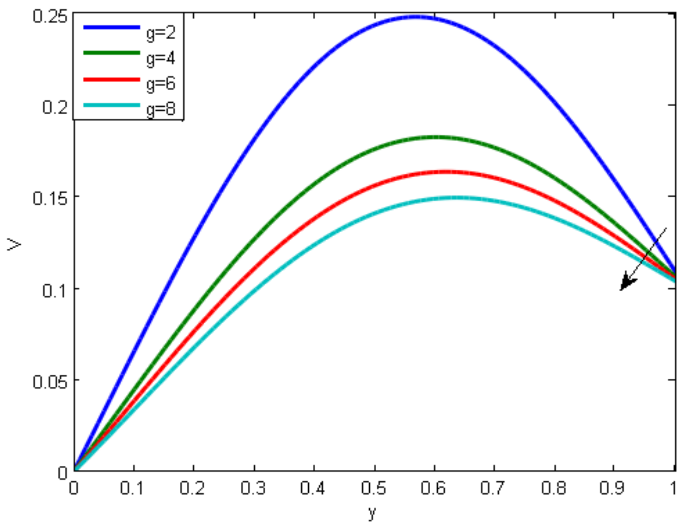

Fig 9: Graph of particle velocity V for Gr,[ϕ = 1,M=2,Pe=2,Re=1,K=1,m=1]

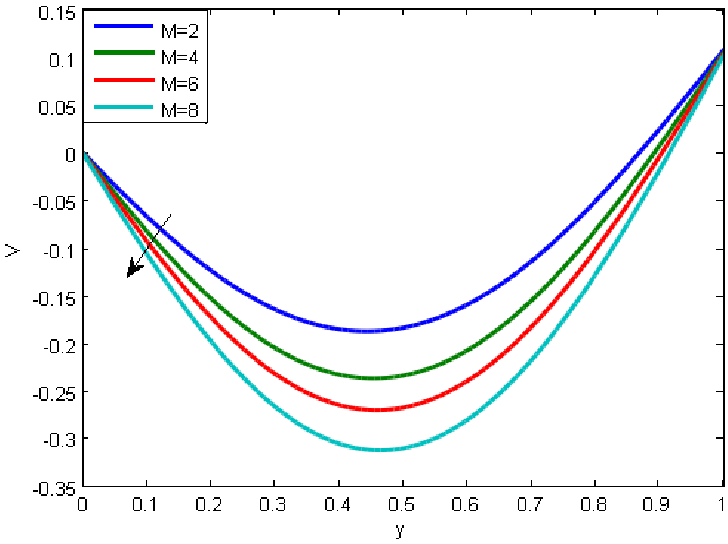

Fig 10: Particle velocity profile V for M, [ϕ =1,Gr = 2, Pe=2,Re=1,K=1,m=1]

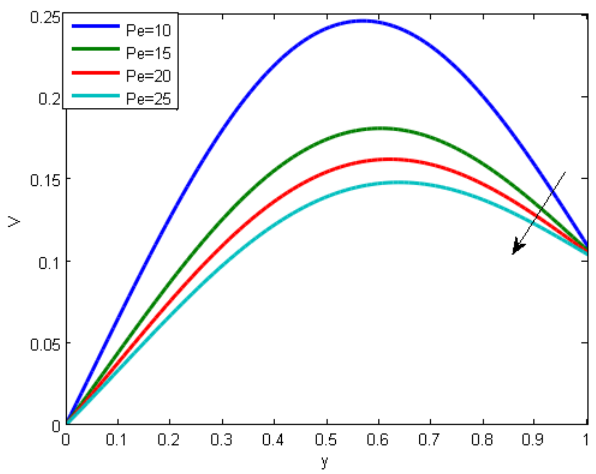

Fig 11: Stetch of particle velocity V for Pe, [ϕ =1,Gr = 2, M=2,Re=1,K=1,m=1]

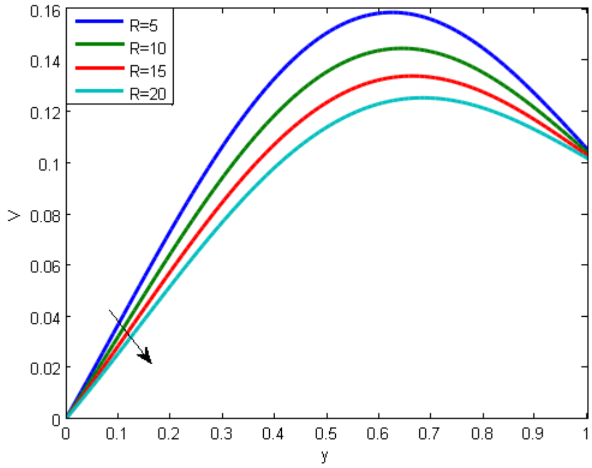

Fig 12: Profile of particle velocity V for R, [ϕ =1,Gr = 2, Pe=2,M=2,K=1,m=1]

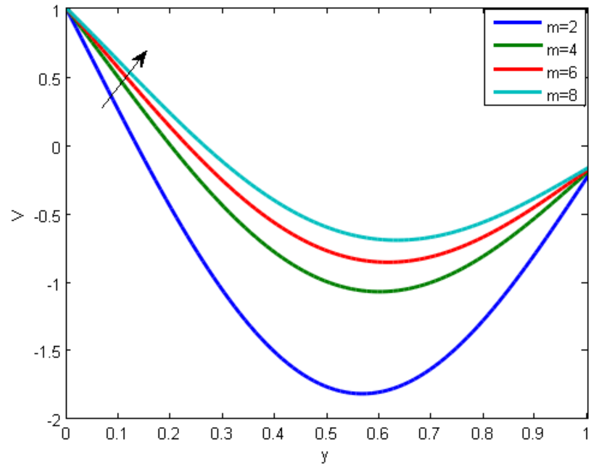

Fig 13: Particle velocity profile V for m, [ϕ =1,Gr = 2, Pe=2,M=2,K=1,m=1]

Fig 14 - 19 shows the temperature profile.It is clearly noted that,by increasing numerical values of the parameters ϕ , N,Gr and m in the flow, an accelerating temperature profile is observed .

The temperature graph of Pe and R is depicted in fig 16 and 17. From the graph the temperature retardation is noted for increasing values of Pe and R.

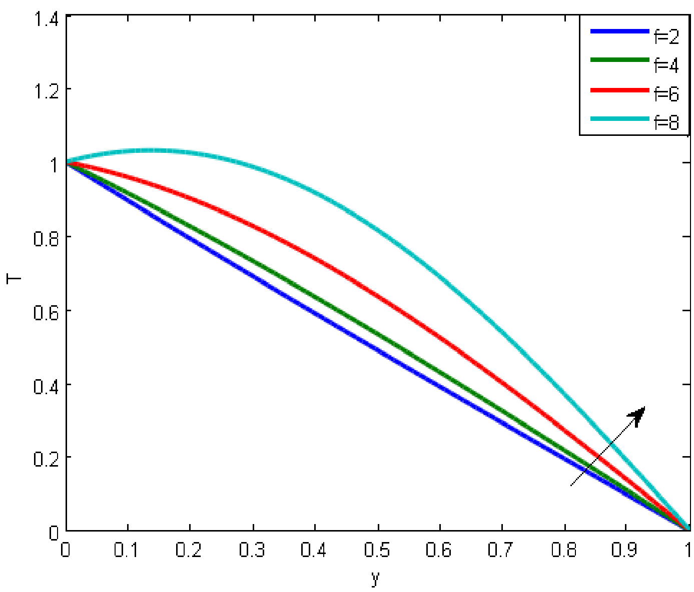

Fig 14: Temperature profile T for ϕ, [R=1,Gr = 2, Pe=2,M=2,K=1,m=1,N=1]

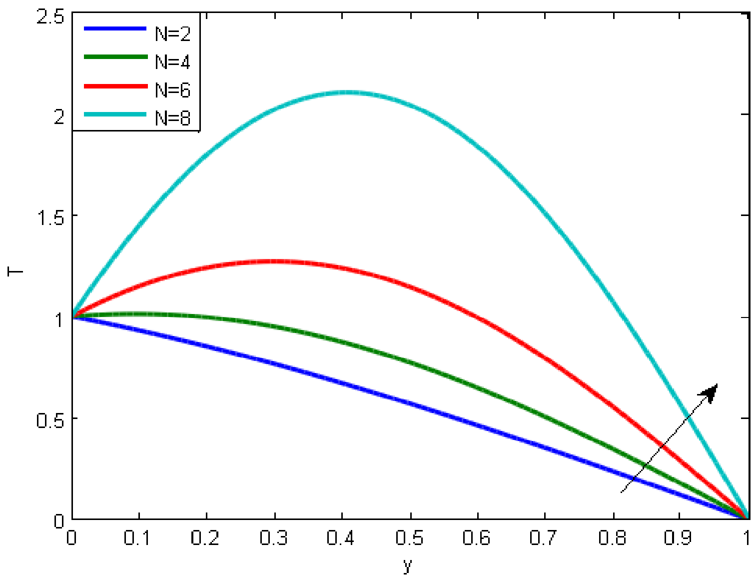

Fig 15: Temperature profile T for N, [ϕ =1,Gr = 2, Pe=2,M=2,K=1,m=1R=1]

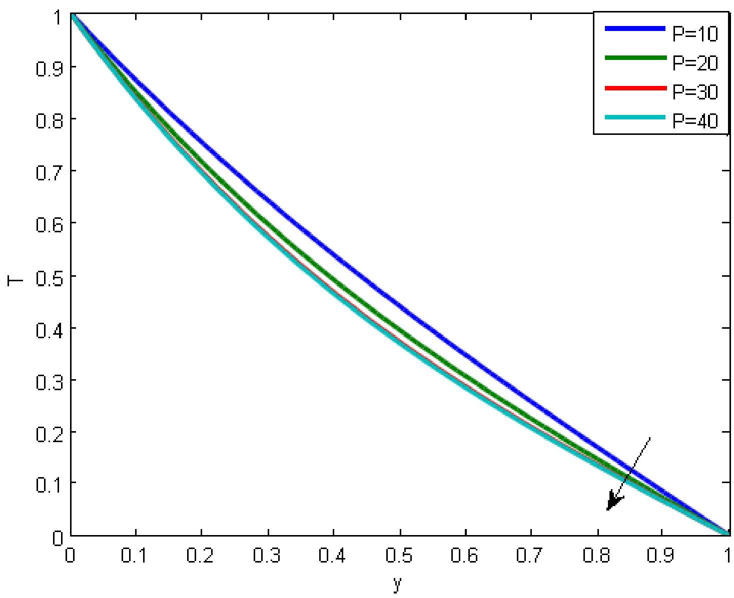

Fig 16: Temperature profile T for Pe, [R=1,Gr = 2, ϕ =1, M=2,K=1,m=1,N=1]

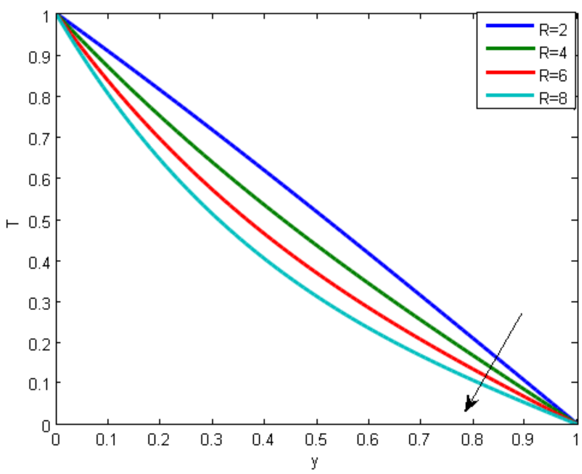

Fig 17: Temperature profile T for R, [Pe=2,Gr = 2, ϕ =1, M=2,K=1,m=1,N=1]

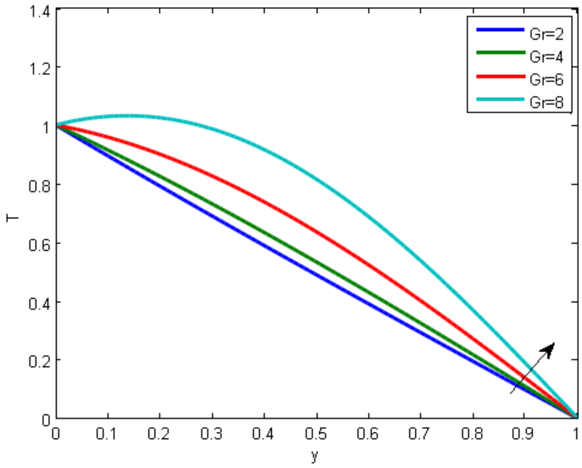

Fig 18: Temperature profile T for Gr, [Pe=2,R = 1, ϕ =1, M=2,K=1,m=1,N=1]

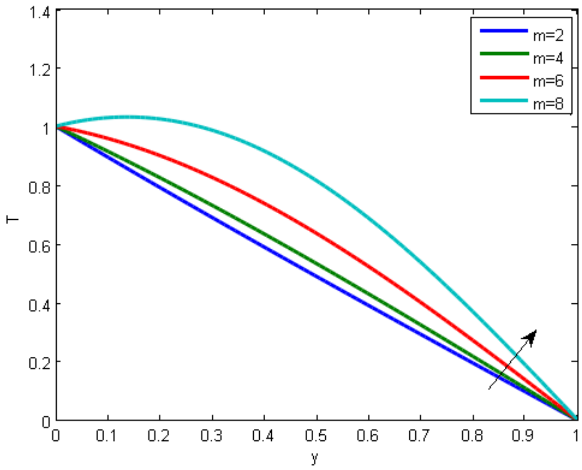

Fig 19: Temperature profile T for m, [Pe=2,R = 1, ϕ =1, M=2,K=1,Gr=2,N=1]

4. Conclusion

The present work eillustrates about the unsteadymagnetohydrodynamic fluctuating flow of viscoelastic fluid bounded by horizontal plates.It is considered that the free stream oscillatory , electrically conducting incompressible with heat radiation has been considered.It is interesting to conclude that the on upgrading Gr,Φ and K parameters,the velocity profile increasing.Considering the velocity profile for dusty particles ,the profile degrades by increasing R,Pe,M,Φ andϒ parameters.Whereas all the parameters involved in the flow shows an increasing temperature profile .

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