

Solving Coupled Fractional Differential Equations Using Differential Transform Method.

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Abstract : This paper presents the solution of coupled equations which are of fractional order using differential transform method. In this paper we extend the scope of differential transform method to system of fractional differential equations so that we get the analytical solutions. The coupled fractional differential equations of a physical system, namely, coupled fractional oscillator with some applications is given via differential transform method. Here we introduce the solution of coupled oscillation of equal fractional order which can be enhanced to unequal fractional order.

Key words: differential transform method, coupled fractional differential equations, coupled oscillation.

2010 subject classification: 34C15, 35R11

1. Introduction

In this paper we have considered system FDE's

$$D_*^{\alpha_1} x_1(t) = f_1(t_1, x_1, x_2 \dots x_n)$$

$$D_*^{\alpha_2} x_2(t) = f_2(t_1, x_1, x_2 \dots x_n) \dots$$

$$D_*^{\alpha_n} x_n(t) = f_n(t_1, x_1, x_2 \dots x_n) \quad (1)$$

Where $D_*^{\alpha_i}$ is derivative of x_i in the sense of Caputo $0 < \alpha_i < 1$ subject to initial conditions.

The differential equations of fractional order have been studied from many years due to their tremendous number applications not only in fluid mechanics and physics but in biology and engineering as well. Application of FDE of nonlinear in nature has been concluded in [1-6].

The DTM first came into existence in engineering domain which has been fully evaluated in [7]. There is a comparison between DTM and ADM for solving FDE as discussed in [8].

The Riemann- Liouville fractional integration of order β is given by

$$J_{x_0}^x f(x) = \frac{1}{\Gamma(\beta)} \int_{x_0}^x (x-t)^{\beta-1} f(t) dt, \beta > 0, x > 0 \quad (2)$$

Given below are the next two equations of Riemann- Liouville and Caputo fractional derivative of order β respectively:

$$D_{x_0}^\beta f(x) = \frac{d^p}{dx^p} [J^{p-\beta} f(x)] \quad (3)$$

$$D_{x_0}^\beta f(x) = J^{p-\beta} \left[\frac{d^p}{dx^p} f(x) \right] \quad (4)$$

Where $p - 1 < \beta \leq p$ and $p \in N$

In the formation of problem on initial and boundary conditions, Caputo fractional derivative has been chosen, and the two operators coincide for homogeneous initial conditions.

2. Fractional differential transform method.

The initial development of FDTM in [8] is as follows:

The fractional differentiation of Riemann- Liouville sense is defined by

$$D_{x_0}^\gamma f(x) = \frac{1}{\Gamma(p-r)} \frac{d^p}{dx^p} \left[\int_{x_0}^x \frac{f(t)}{(x-t)^{1+r-p}} dt \right], \tag{5}$$

For $p - 1 \leq \gamma \leq p$ and $p \in \mathbb{Z}^+, x > x_0$. We can express it in power series as follows:

$$f(x) = \sum_{k=0}^{\infty} F(k)(x - x_0)^{\frac{k}{\beta}} \tag{6}$$

Where β is order of the fraction and $F(k)$ is fractional differential transform of $f(x)$.

Avoiding initial and boundary condition of fractional order, we define derivative in Caputo sense. Relation between operators of Riemann- Liouville and Caputo is given as

$$D_{*x_0}^\gamma f(x) = D_{x_0}^\gamma [f(x) - \sum_{k=0}^{p-1} \frac{1}{k!} (x - x_0)^k f^{(k)}(x_0)] \tag{7}$$

Setting $f(x) = f(x) - \sum_{k=0}^{p-1} \frac{1}{k!} (x - x_0)^k f^{(k)}(x_0)$ in Eq. 5 and using Eq. 7 we obtain Caputo sense as

$$D_{*x_0}^\gamma f(x) = \frac{1}{\Gamma(p-r)} \frac{d^p}{dx^p} \left[\int_{x_0}^x \frac{f(t) - \sum_{k=0}^{p-1} \frac{1}{k!} (t - x_0)^k f^{(k)}(x_0)}{(x-t)^{1+r-p}} dt \right]$$

Applying initial conditions to integer order derivatives,

$$F(k) = \begin{cases} \text{if } \frac{k}{\beta} \in \mathbb{Z}^+, \frac{1}{(\frac{k}{\beta})!} \left[\frac{d^{\frac{k}{\beta}} f(x)}{dx^{\frac{k}{\beta}}} \right]_{x=x_0} & (8) \\ \text{if } \frac{k}{\beta} \notin \mathbb{Z}^+ & 0 \end{cases}$$

We can obtain below mentioned theorems from 5 and 6, for proof and details see [8].

Theorem 1. If $h(x) = u(x) \pm s(x)$ implies $H(k) = U(k) \pm S(k)$

Theorem 2. If $h(x) = u(x)s(x)$, implies $H(k) = \sum_{l=0}^k U(l)S(k-l)$

Theorem 3. If $h(x) = u_1(x)u_2(x) \dots \dots \dots u_n(x)$ then

$$H(k) = \sum_{k_{n-1}=0}^k \sum_{k_{n-2}=0}^{k_{n-1}} \sum_{k_{n-3}=0}^{k_{n-2}} \sum_{k_{n-2}=0}^{k_{n-1}} \dots \dots \dots \sum_{k_2=0}^{k_3} \sum_{k_1=0}^{k_2} U_1(k_1)U_2(k_2 - k_1) \dots U_n(k - k_{n-1})$$

Theorem 4. If $h(x) = (x - x_0)^m$ then $H(k) = \varphi(k - \beta m)$

$$\varphi(k) = \begin{cases} 1 & \text{if } k = 0 \\ 0 & \text{if } k \neq 0 \end{cases}$$

Theorem 5. If $h(x) = D_{x_0}^\gamma [r(x)]$ implies $(k) = \frac{\Gamma(r+1+\frac{k}{\beta})}{\Gamma(1+\frac{k}{\beta})} U(k + \beta r)$.

Theorem 6. In most general form, fractional derivatives can be produced from

$h(x) = \frac{d^{r_1}}{dx^{r_1}} [u_1(x)] \frac{d^{r_2}}{dx^{r_2}} [u_2(x)] \dots \frac{d^{r_{n-1}}}{dx^{r_{n-1}}} [u_{n-1}(x)] \frac{d^{r_n}}{dx^{r_n}} [u_n(x)]$, implies

$$H(k) = \sum_{k_{n-1}=0}^k \sum_{k_{n-2}=0}^{k_{n-1}} \sum_{k_{n-3}=0}^{k_{n-2}} \sum_{k_{n-2}=0}^{k_{n-1}} \dots \sum_{k_2=0}^{k_3} \sum_{k_1=0}^{k_2} \frac{\Gamma\left(r_1 + 1 + \frac{k_1}{\beta}\right) \Gamma\left(r_2 + 1 + \frac{k_2 - k_1}{\beta}\right)}{\Gamma\left(1 + \frac{k_1}{\beta}\right) \Gamma\left(1 + \frac{k_2 - k_1}{\beta}\right)} \dots \frac{\Gamma\left(r_2 + 1 + \frac{k - k_n}{\beta}\right)}{\Gamma\left(1 + \frac{k - k_n}{\beta}\right)}$$

$$U_1(k_1 + \beta r_1) U_2(k_2 - k_1 + \beta r_2) \dots U_n(k - k_{n-1} + \beta r_n); \beta r_j \in z^+ \text{ for } j = 1, 2, \dots, n$$

3. Coupled fractional differential equation

We introduce the DTM to solve linear inhomogeneous coupled fractional differential equations. Our discussion is restricted to only linear non-homogenous ODE's of fractional order. As given in [9] Coupled oscillators are solved using various techniques. Here we solve the same oscillator equation with DTM to make it more convenient.

To proceed with, linear-coupled oscillator system is given by

$$D^{\gamma_1} x_1(t) = -\omega_2 x_1(t) + \bar{\omega}_2 \{x_2(t) - x_1(t)\},$$

$$D^{\gamma_2} x_2(t) = -\omega_2 x_2(t) + \bar{\omega}_2 \{x_1(t) - x_2(t)\}$$

Here we use the substitution as

$$-\omega_2 + \bar{\omega}_2 = \mu \ \& \ \bar{\omega}_2 = \epsilon$$

The above equations become

$$D^{\gamma_1} x_1(t) = -\mu x_1(t) + \epsilon x_2(t)$$

$$D^{\gamma_2} x_2(t) = -\mu x_2(t) + \epsilon x_1(t) \tag{9}$$

Applying parameters; $\omega = 1, \bar{\omega} = 0.5, x_1(0) = 0, x_2(0) = 1$

$$D^{\gamma_1} x_1(0) = 0$$

$$D^{\gamma_2} x_2(0) = 0.1$$

System (1.9) can be changed by using Theorems 1 to 5

$$X_1(k + \gamma_1 \beta_1) = \frac{\Gamma\left(1 + \frac{k}{\beta_1}\right)}{\Gamma\left(\gamma_1 + 1 + \frac{k}{\beta_1}\right)} [-\mu X_1(k) + \epsilon X_2(k)]$$

$$X_2(k + \gamma_2 \beta_2) = \frac{\Gamma\left(1 + \frac{k}{\beta_2}\right)}{\Gamma\left(\gamma_2 + 1 + \frac{k}{\beta_2}\right)} [\epsilon X_1(k) - \mu X_2(k)] \tag{10}$$

$$X_1(k) = 0 \quad \text{for } k = 1, 2, 3 \dots \gamma_1 \beta_1 - 1$$

$$X_2(k) = 0 \quad \text{for } k = 1, 2, 3 \dots \gamma_2 \beta_2 - 1$$

$$X_2(0) = 1 \tag{11}$$

For $\gamma_1 = 1$ and $\gamma_2 = 1; \beta_1 = 1$ and $\beta_2 = 1$

$$x_1(0) = 0 \quad ; \quad x_2(0) = 1$$

$$x_1(1) = \epsilon \quad ; \quad x_2(1) = -\mu$$

$$x_1\left(\frac{1}{2}\right) = \frac{1}{2}[-\epsilon\mu - \mu\epsilon] = -\mu\epsilon$$

$$x_2(1) = \frac{1}{2}[\mu^2 + \epsilon^2]$$

Then following series can be obtained for $x_1(t)$ and $x_2(t)$, respectively.

$$x_1(t) = \epsilon t - \mu t^2 + \dots$$

$$x_2(t) = -1 - \mu t + \frac{1}{2}[\mu^2 + \epsilon^2]t^2 + \dots$$

For $\gamma_1 = 1.4$ and $\gamma_2 = 1.7$ we have

$$X_1(t + 14) = \frac{\Gamma(1 + k/10)}{\Gamma(14 + 1 + k/10)} [-\mu X_1(k) + \epsilon X_2(k)]$$

$$X_2(t + 17) = \frac{\Gamma(1 + k/10)}{\Gamma(1.7 + 1 + k/10)} [X_1(k) - \mu X_2(k)]$$

For k up to 50 and then using in above equation we get the below mentioned series for $x_1(t)$ and $x_2(t)$, respectively.

$$x_1(t) = \frac{\epsilon}{\Gamma(2.4)} t^{\frac{14}{10}} - \frac{\epsilon\mu}{\Gamma(3.8)} t^{\frac{28}{10}} - \frac{\epsilon\mu}{\Gamma(2.7)} t^{\frac{31}{10}} + \frac{\Gamma(4.1)}{\Gamma(5.5)} \left[\frac{\epsilon\mu^2}{\Gamma(2.7)} + \frac{\epsilon^3}{\Gamma(4.1)} \right] t^{\frac{45}{10}} + \frac{\epsilon\mu^2}{\Gamma(4.4)} t^{\frac{48}{10}} + \dots$$

$$x_2(t) = 1 - \frac{\mu}{\Gamma(2.7)} t^{\frac{17}{10}} + \frac{\epsilon^2}{\Gamma(4.1)} t^{\frac{31}{10}} + \frac{\mu^2}{\Gamma(4.4)} t^{\frac{34}{10}} - \frac{\mu\epsilon^2}{\Gamma(5.5)} t^{\frac{45}{10}} + \frac{\Gamma(4.1)}{\Gamma(5.8)} \left[\frac{\mu\epsilon^2}{\Gamma(2.7)} + \frac{\mu\epsilon^2}{\Gamma(4.1)} \right] t^{\frac{48}{10}} + \dots$$

4. Conclusion

The present work assures the applicability of solving coupled fractional differential equation by DTM. Earlier it was done with various numerical methods like Laplace transformation method, adjoint method etc. The much greater advantage of this method of solving coupled fractional differential equation over adomain decomposition method is that we do not get adomain polynomials. This method is a reliable one and full of techniques and much easier in solving fractional differential equations of coupled form both linear and nonlinear.

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