Dynamics Of the Polynomial Differential Systems

Feneniche Fatima, Rezaoui Med Salem[†],

1,2 University of Sciences (USTHB), Faculty of Mathematics, B.P 32, El-Alia,16111, Bab-Ezzouar, Algiers, Laboratory of Algebra and Number Theory (LATN)

Fatima Feneniche – received a Master degree in Dynamics Mathematics from Science and Technology Houari Boumediene University (USTHB), Bab-Ezzouar, Algeria,

^bMed-Salem Rezaoui –is working as Professor in the Department of Mathematics, Science and Technology Houari Boumediene University (USTHB), Bab-Ezzouar, Algeria. His research area includes Differential Systems and Applications.

Article History: Received 3 June 2022; Accepted 14 January 2023; Published 3 September 2023

Abstract: In this paper, we study the bifurcation of limit cycles for the following Lienard systems

x' = y, $y' = -f_m(x)y - g_n(x)$,

where, $f_m(x)$ and $g_n(x)$ respectively are polynomials of degree m and n, $g_n(0) = 0$. We prove that, if m = 5 and $g_n(x) = x$, there are always Lienard systems of the above form as they have a limit cycle.

Keywords: Limit cycles, The bifurcation set.

1. Introduction

In general, there are several free parameters. By using a method introduced in a previous paper, we obtain a sequence of algebraic a proximations to the bifurcationsets, in the parameter space.

Each algebraic approximation represents an exact lower bound to the bifurcation set.

The method is perturbative. So, it is not necessary to have a small or a large parameter in order to obtain these results.

We consider the following problem

$$\begin{cases} \frac{dx}{dt} = y\\ \frac{dy}{dt} = \varepsilon(1 - x^2)y - x. \end{cases}$$

(See [1]).

Liénard system

In 1926, German Van Der Pol proposed the differential equation

 $\ddot{x} + \varepsilon (x^2 - 1)\dot{x} + x = 0, \qquad \varepsilon > 0.$ (1)

Can the Lienard system with m > 5 and m + 1 < n < 2m have an algebric limit cycle ?

In this part, by developing the main ideas we prove the next results which give a positive solution to the problem opened above.

2. MAIN RESULT

The first step We choose m = 5 and n = 1 then $f_m(x) = a_0 x^5 + a_1$, $g_n(x) = x$.

The system becomes.

[†] Fatima Feneniche. Email addresses: 1 ffeneniche@usthb.dz & fatima.feneniche@gmail.com Med-Salem REZAOUI. E-mail addresses: 2mrezaoui@usthb.dz & s_rezaoui@yahoo.fr

$$\begin{cases} \dot{x} = y \\ \dot{y} = -(a_0 x^5 + a_1)y - x. \end{cases}$$

First approximation the general solution is

 $x = C \cos(t - \alpha)$:

(See [2])

Where C and α are arbitrary, we study how the presence of the term with ϵ affects on the solution

 $x = C \cos t$,

Being given what should be a first reasonable approximation, we try

 $x = Acos(\omega t),$

Where ω is a constant which is close to 1 (See [2]).

$$\begin{cases} \dot{x} = -A\omega\sin(\omega t) \\ \ddot{x} = -A\omega^2\cos(\omega t). \end{cases}$$

by replacing this in the coming differential equation

$$\ddot{x} + (a_0 x^5 + a_1) \dot{x} + x = 0, \qquad (2)$$

We find that
$$-Aw^2 \cos(\omega t) - (a_0 A^5 \cos(\omega t)^5 + a_1)(Aw \sin(\omega t)) + A\cos(\omega t) = 0$$

 $A(-w^2)\cos(\omega t) = a_0 A^6 w \cos(\omega t)^5 \sin(\omega t) + a_1 Aw \sin(\omega t),$
 $= Aw(a_0 A^5 \cos(\omega t)^5 \sin(\omega t) + a_1 A \sin(\omega t))$.

According to the simplification, we find the second member equal at :

$$= \operatorname{Aw}\left(a_{0}\frac{A^{5}}{32}\sin(6\omega t) + \frac{3}{16}a_{0}A^{5}\sin(4\omega t) + \frac{15}{32}a_{0}A^{5}\sin(2\omega t) + a_{1}\sin(\omega t)\right).$$

This equation can be satisfactory for all t only if the coefficients of the different terms sinusoidal disappear The term cos(wt) diappear if w=1 and the coefficient of sin(wt) is zero if we take

$$w=1, a_1=0.$$

We see that the choice of A is arbitrary that signifies that the system doesn't admit an isolated closed curve as limited cycle (See [9]). So, we search for another more efficient method :

Passage in coordinated polar We take :

$$\begin{aligned} x &= r\cos(\omega t) \\ y &= r\sin(\omega t). \end{aligned}$$

According to the study,was :

$$\begin{cases} \dot{r} = (x\dot{x} + y\dot{y})/r\\ \dot{\theta} = (x\dot{y} - y\dot{x})/r^2. \end{cases}$$

(See [11,13]) We have :

$$x\dot{x} + y\dot{y} = xy - (a_0x^5 + a_1)y^2 - xy,$$

= $-(a_0x^5 + a_1)y^2,$
= $-r^2(a_0r^5\cos(\theta)^5 + a_1)\sin(\theta)^2.$

Then

$$\begin{split} \dot{r} &= -r(a_0 r^5 \cos \theta^5 + a_1) \sin \theta^2, \\ &= -r(a_0 r^5 \cos \theta^5 \sin \theta^2 + a_1 \sin \theta^2), \\ &= -r(a_0 r^5 \left(\frac{1}{16} \cos 5\theta + \frac{5}{16} \cos 3\theta + \frac{5}{8} \cos \theta\right) \sin \theta^2 + a_1 \frac{1 - \cos 2\theta}{2}), \\ \dot{r} &= \frac{3}{64} a_0 r^6 \cos 5\theta + \frac{1}{64} a_0 r^6 \cos 7\theta + a_1 \frac{r^2}{2} \cos 2\theta - \frac{5}{64} a_0 r^6 \cos \theta - a_1 \frac{r^2}{2} + \frac{1}{64} a_0 r^6 \cos 3\theta. \end{split}$$

We have :

$$\begin{aligned} x\dot{y} - y\dot{x} &= -(a_0x^5 + a_1)xy - x^2 - y^2, \\ &= -(a_0x^5 + a_1)xy - (x^2 + y^2), \\ &= -r^2((a_0r^5\cos 5\theta + a_1)\cos \theta - \sin \theta + 1) \end{aligned}$$

We have :

$$\dot{\theta} = -((a_0 r^5 \cos 5\theta + a_1) \cos \theta \quad \sin \theta + 1),$$

$$\dot{\theta} = -(a_0 r^5 \frac{1}{16} \cos 5\theta + \frac{5}{16} \cos 3\theta + \frac{5}{8} \cos \theta + a_1) \frac{\sin 2\theta}{2} + 1),$$

$$\dot{\theta} = a_0 r^5 \left(\frac{1}{32}\cos 5\theta \sin 2\theta - \frac{5}{32}\cos 3\theta \sin 2\theta - \frac{5}{16}\cos \theta \sin 2\theta - a_1 \frac{\sin 2\theta}{2} - 1\right),$$

$$\dot{\theta} = -\frac{1}{64}a_0 r^5 \sin 7\theta - \frac{5}{64}a_0 r^5 \sin 5\theta - \frac{5}{64}a_0 r^5 \sin 3\theta + \frac{15}{64}a_0 r^5 \sin \theta - \frac{a_0 r^5 a_1}{2}\sin 2\theta - 1$$

So the system becomes :

$$\begin{cases} \dot{r} = \frac{3}{64}a_0r^6\cos 5\theta + \frac{1}{64}a_0r^6\cos 7\theta + a_1\frac{r^2}{2}\cos 2\theta - \frac{5}{64}a_0r^6\cos \theta - a_1\frac{r^2}{2} + \frac{1}{64}a_0r^6\cos 3\theta. \\ \dot{\theta} = -\frac{1}{64}a_0r^5\sin 7\theta - \frac{5}{64}a_0r^5\sin 5\theta - \frac{5}{64}a_0r^5\sin 3\theta + \frac{15}{64}a_0r^5\sin \theta - \frac{a_0r^5a_1}{2}\sin 2\theta - 1. \end{cases}$$

We estimate the conditions starting from the existence of the limited cycle (th Bedixon)(See [9,6,3,10])

, so that the hypothesis doesn't change $a_0; a_1 \neq 0$.

Perturbation of the system We take the equation:

$$\ddot{x} + (a_0 x^5 + a_1) \dot{x} + x = 0,$$

Development in the neighbourhood limit (0, 0). (See [9])

For the value of $\varepsilon = 0$, we find the exact solution.

The wanted idea is to extend the solution in power serial ε :

$$x(t) = x_0(t) + x_1(t)\varepsilon + x_2(t)\varepsilon^2 + 0(\varepsilon^3).$$
 (3)

So:

$$=a_0(x_0^{5}(t) + 5x_0^{4}(t)x_1(t)\varepsilon + 10x_0^{3}(t)x_1^{2}(t)\varepsilon^{2} + 5x_2(t)x_0^{4}(t)\varepsilon^{2}) + a_1.$$

Is replaced in equation:

$$\begin{aligned} \ddot{x}(t) + x(t) &= (\ddot{x}_0(t) + x_0(t)) + (\ddot{x}_1(t) + x_1(t))\varepsilon + (\ddot{x}_2(t) + x_2(t))\varepsilon^2 + 0(\varepsilon^3). \\ a_1\dot{x}(t) &= a_1\dot{x}_0(t) + a_1\dot{x}_1(t)\varepsilon + a_2\dot{x}_2(t)\varepsilon^2 + 0(\varepsilon^3). \\ a_0x^5(t)\dot{x}(t) &= a_0x_0^{-5}(t)\dot{x}_0(t) + 5a_0x_0^4(t)x_1(t)\dot{x}_0(t)\varepsilon + 10a_0x_0^{-3}(t)x_1^{-2}(t)\dot{x}_0(t)\varepsilon^2 \\ &+ 5a_0x_0^4(t)x_2(t)\dot{x}_0(t)\varepsilon^2 + \\ a_0x_0^{-5}(t)\dot{x}_1(t)\varepsilon + 5a_0x_0^4(t)x_1(t)\dot{x}_1(t)\varepsilon^2 + a_0x_0^{-5}(t)\dot{x}_2(t)\varepsilon^2. \end{aligned}$$

We find

$$\begin{aligned} (x_0 + \ddot{x_0} + a_1 \dot{x_0} + a_0 x_0^5 \dot{x_0}) + (x_1 + \ddot{x_1} + a_1 \dot{x_1} + a_0 x_0^5 \dot{x_1} + 5a_0 x_0^4 x_1 \dot{x_0})\varepsilon \\ &+ (x_2 + \ddot{x_2} + a_1 \dot{x_2} + \ddot{10} a_0 x_0^3 x_1^2 \dot{x_0} + 5a_0 x_0^4 x_2 \dot{x_0} + 5a_0 x_0^4 x_1 x_1 + a_0 x_0^5 \dot{x_2})\varepsilon^2 + 0(\varepsilon^3) \\ &= 0 \end{aligned}$$

We want find a solution that is valid for all little values of ε . We cancel each of the coefficient $of \varepsilon^n$.

For $n = 0, 1, 2, \dots$

For n = 0; 1 and 2, we obtained :

$$\begin{aligned} x_{0} + \ddot{x_{0}} + a_{1}\dot{x_{0}} + a_{0}x_{0}{}^{5}\dot{x_{0}} &= 0 \end{aligned} (4) \\ x_{1} + \ddot{x_{1}} + a_{1}\dot{x_{1}} + a_{0}x_{0}{}^{5}\dot{x_{1}} + 5a_{0}x_{0}{}^{4}x_{1}\dot{x_{0}} &= 0 \end{aligned} (5) \\ x_{2} + \ddot{x_{2}} + a_{1}\dot{x_{2}} + \ddot{10}a_{0}x_{0}{}^{3}x_{1}{}^{2} \dot{x_{0}} + 5a_{0}x_{0}{}^{4}x_{2}\dot{x_{0}} + 5a_{0}x_{0}{}^{4}x_{1}x_{1} + a_{0}x_{0}{}^{5}\dot{x_{2}} \\ &= 0 \end{aligned} (6)$$

As the equation is autonomous we can choose the instant for corresponding t = 0 to any point of the limited cycle.

Thus, we can choose the initial condition $\dot{x}(0)=0$ with out losing genralty from developing (5).

we obtain the initial conditions :

$$\dot{x}_0(0) = \dot{x}_1(0) = \dot{x}_2(0) = \dots = 0 \quad (7)$$

EQS (6) and (9) given
 $x_{0h}(t) = \alpha \cos(\theta)$, homogeneous solution. (8)

Where α is yet determined.

we find the particular solution equal at :

$$x_{0p}(t) = A_0 \cos(6\theta) + B_0 \cos(4\theta) + R_0 \cos(2\theta) + M_0 \cos(\theta).$$

We need of all terms (are not periodic) A_0 , B_0 , R_0 , M_0 to be removed .by substituting (10) in (7) and by using trigonometric identities, we obtain :

We need of all terms (are not periodic) A_0 ; B_0 ; R_0 ; M_0 to be removed .

by substituting (10) in (7) and by using trigonometric identities.

We obtain :

$$\ddot{x_1} + x_1 = -(a_1 + a_0 \alpha^5 \cos \theta^5) x_1 + 5a_0 \alpha^5 \cos \theta^4 \sin \theta x_1$$
(9)
$$\ddot{x_1} + (a_1 + a_0 \alpha^5 \cos \theta^5) \dot{x_1} + (1 - 5a_0 \alpha^5 \cos \theta^4 \sin \theta) x_1 =$$

(10) as a result, this is a

0

differential equation of second order with variable coefficients. This equation is ,relatively, easy to solve in the general case. Be the equation

Supposing that the function x_1 that satisfies the differential equation $x_1 = \exp(k t)$ where k can be a complex number.

So we have

or

$$k^{2}exp^{kt} + (a_{1} + a_{0}\alpha^{5}\cos\theta^{5})kexp^{kt} + (1 - 5\alpha^{5}\cos\theta^{4}\sin\theta)exp^{kt} = 0,$$
$$k^{2} + (a_{1} + a_{0}\alpha^{5}\cos\theta^{5})k + (1 - 5\alpha^{5}\cos\theta^{4}\sin\theta) = 0,$$

this last relation is the quadratic equation auxiliary of the differential equation (polynomial charac-

teristic)(See [9]).

It has to be solutions / roots that we will notice in the general case : k_1, k_2 .

Immediately comes that :

$$k_{1,2} = \frac{-(a_1 + a_0\alpha^5\cos\theta^5) \pm \sqrt{(a_1 + a_0\alpha^5\cos\theta^5)^2 - 4(1 - 5\alpha^5\cos\theta^4\sin\theta)}}{2}$$

So

$$x_1 = s_1 exp^{k_1 t} + s_2 exp^{k_2 t}$$

If we take $a_0 = 0$ or $\alpha = 0$ we find that x_1 is a periodic solution, but it's contradictory to the hypothesis or to the trivial solution x = 0.

Bifurcation system We have the following system:

$$\begin{cases} \dot{x} = y \\ \dot{y} = -(a_0 x^5 + a_1)y - x. \end{cases}$$

We search equilibrium points for this system :

$$\begin{cases} \dot{x} = 0\\ \dot{y} = 0 \end{cases}$$

The system accepts one equilibrium point which is : (0; 0).

After that, we look for the Jacobean:

$$\mathbf{J}(\mathbf{x},\mathbf{y}) = \begin{pmatrix} 0 & & 1 \\ -(5a_0x^5y) - 1 & & -(a_0x^5 + a_1) \end{pmatrix}$$

So, at the equilibrium point:

$$\mathbf{J}(0,0) = \begin{pmatrix} 0 & 1\\ -1 & -a_1 \end{pmatrix}$$

We search for the polynomial characteristic :

$$\lambda^2 + a_1 \lambda + 1 = 0.$$
 $\delta = a_1^2 - 4$

We find three cases δ :

If $\delta = 0$ then $a_1 = \pm 2$ thus the equilibrium point is a node.

If $\delta > 0$ then $a_1 > 2$, thus λ it accepts two values of different signs, and the equilibrium point is a saddle point.

The interesting case is $\delta < 0$, then the value of λ is complex and in this case, we have :

$$\lambda = -a_1 \pm \sqrt{a_1^2 - 4}.$$

If $a_1 = 0$ the equilibrium points is central in the linear case for the non linear part it's a passage of polar coordinated because linear is topologically equivalent with the non linear but the centre can be a limited cycle.

If $0 < a_1 < 2$ the equilibrium point is an unstable focus (positive real part).

If $0 < a_1$ the equilibrium point is a steady focus (negative real part).

In the non linear case, we can use lyaponov's criterium :

The system linear predicts centres when the setting is equal to 0.

In the non linear system, we can see that those centres aren't in fact conserved. To determine the steadiness of the origin, we consider positive defined function v(x, y). I take some examples of v(x, y) but they aren't efficient.

The second step

We choose :

$$f_m(x) = a_0 x^5 + a_1 x^4 + a_2,$$
 $g_n(x) = x.$

The system becomes :

$$\begin{cases} \dot{x} = y \\ \dot{y} = -(a_0 x^5 + a_1 x^4 + a_2)y - x \end{cases}$$

In the first approximation, the general solution is :

$$x = C \cos(t - \alpha).$$

Where *C* and α are arbitrary.

We study how the presence of the term with ε affects on the solution :

$$x = C cost.$$

By giving what can be a first reasonable approximation, we try

$$x = A \cos(\omega t)$$
.

Where w is a constant close to 1.

$$\begin{cases} \dot{x} = -A \, w \sin wt \\ \ddot{x} = -Aw^2 \cos wt \end{cases}$$

We replace in the following differential equation

$$\ddot{x} + (a_0 x^5 + a_1 x^4 + a_2) \dot{x} + x = 0.$$
(12)

We find that :

$$\begin{aligned} -Aw^{2}\cos wt - (a_{0}A^{5}\cos wt^{5} + a_{1}A^{4}\cos wt^{4} + a_{2})(Aw\sin wt + A\cos wt) &= 0, \\ A(1 - w^{2})\cos wt &= a_{0}A^{6}w\cos wt^{5}\sin wt + a_{1}A^{5}w\cos wt^{4}\sin wt + a_{2}Aw\sin wt, \\ &= Aw(a_{0}A^{5}\cos wt^{5} + a_{1}A^{4}\cos wt^{4} + a_{2})\sin wt. \end{aligned}$$

After the simplification, we find the second member equal at :

$$= a_0 \frac{A^6 w}{32} \sin 6wt$$

+ $a_1 \frac{A^5 w}{16} \sin 5wt + a_0 \frac{A^6 w}{8} \sin 4wt$
+ $3a_1 \frac{A^5 w}{16} \sin 3wt + a_0 \frac{A^6 w}{32} \sin 2wt + (a_1 \frac{A^5 w}{8} + a_2 Aw) \sin wt.$

This equation can be satisfied for all t only if the coefficients of the different sinusoidal terms disappear .

The term cos(wt) disappears if w = 1, and the coefficient of $A = \sqrt[4]{-8\frac{a_2}{a_1}}$ is zero as if we take $a_1, a_2 < 0, a_1 \neq 0$.

Thus, we choose :

w = 1, A =
$$\sqrt[4]{-8\frac{a_2}{a_1}}$$
.

This lets yet the term that containing $\sin(6wt)$; $\sin(5wt)$; $\sin(4wt)$; $\sin(3wt)$; $\sin(2wt)$ to disappear because we have fixed w and A.

Thus, we find that the system admits a limited cycle of radius $A = \sqrt[4]{-8\frac{a_2}{a_1}}$ as $a_1, a_2 < 0, a_1 \neq 0$.

3rdStep

We choose

$$f_m(x) = a_0 x^5 + a_1 x^4 + a_2 x^3 + a_3, \qquad g_n(x) = x.$$

The system becomes :

$$\begin{cases} \dot{x} = y \\ \dot{y} = -(a_0 x^5 + a_1 x^4 + a_2 x^3 + a_3)y - x \end{cases}$$

In first approximation, the general solution is

$$x = C \cos(t - \alpha).$$

Where *C* and α are arbitrary.

We study how the presence of the term with ε affects on the solution

$$x = C cost.$$

By giving what can be a first reasonable approximation, we try

$$x = A \cos wt$$
..

Where ω is a constant close to 1.

$$\begin{cases} \dot{x} = -A \, w \sin wt \\ \ddot{x} = -Aw^2 \cos wt \end{cases}$$

We replace in the following differential equation

$$\ddot{x} + (a_0 x^5 + a_1 x^4 + a_2 x^3 + a_3) \dot{x} + x = 0.$$
(13)

We find that :

$$\begin{aligned} -Aw^{2}\cos wt - (a_{0}A^{5}\cos wt^{5} + a_{1}A^{4}\cos wt^{4} + a_{2}A^{3}\cos wt^{3} + a_{3})(Aw\sin wt + A\cos wt) &= 0, \\ A(1 - w^{2})\cos wt &= a_{0}A^{6}w\cos wt^{5}\sin wt + a_{1}A^{5}w\cos wt^{4}\sin wt \\ &+ a_{2}Aw\sin wt + a_{2}A^{4}w\cos wt^{3}\sin wt + a_{3}Aw\sin wt, \end{aligned}$$

$$= Aw (a_0 A^5 \cos wt^5 + a_1 A^4 \cos wt^4 + a_2 A^3 \cos wt^3 + a_3) \sin wt$$

After the simplification, we find the second member equal at :

$$= a_0 \frac{A^6 w}{32} \sin 6wt + a_1 \frac{A^5 w}{16} \sin 5wt + (a_0 \frac{A^6 w}{8} + a_2 \frac{A^4 w}{8}) \sin 4wt + 3a_1 \frac{A^5 w}{16} \sin 3wt + (a_0 \frac{A^6 w}{32} + a_2 \frac{A^4 w}{4}) \sin 2wt + (a_1 \frac{A^5 w}{8} + a_3 Aw) \sin wt.$$

This equation can be satisfied for all t only if the coefficients of the different sinusoidal terms disappear .

The term cos(wt) disappears if w = 1, and the coefficient of sin wt is zero as if we take $A = \sqrt[4]{-8\frac{a_3}{a_1}}$ $a_1, a_3 < 0, a_1 \neq 0.$

This lets yet the term that containing $\sin(6wt)$; $\sin(5wt)$; $\sin(4wt)$; $\sin(3wt)$; $\sin(2wt)$ to disappear because we have fixed w and A.

Thus, we find that the system admits a limited cycle of radius $A = \sqrt[4]{-8\frac{a_3}{a_1}}$ as $a_1, a_3 < 0, a_1 \neq 0$.

4thstep

We choose

$$f_m(x) = a_0 x^5 + a_1 x^4 + a_2 x^3 + a_3 x^2 + a_4, \qquad g_n(x) = x.$$

The system becomes :

$$\begin{cases} \dot{x} = y \\ \dot{y} = -(a_0 x^5 + a_1 x^4 + a_2 x^3 + a_3 x^2 + a_4)y - x \end{cases}$$

In first approximation, the general solution is

$$x = C \cos(t - \alpha).$$

Where *C* and α are arbitrary.

We study how the presence of the term with ε affects on the solution

$$x = C cost.$$

By giving what can be a first reasonable approximation, we try

$$x = A \cos wt$$
..

Where ω is a constant close to 1.

-

$$\begin{cases} \dot{x} = -A \, w \sin wt \\ \ddot{x} = -Aw^2 \cos wt \end{cases}$$

We replace in the following differential equation

$$\ddot{x} + (a_0 x^5 + a_1 x^4 + a_2 x^3 + a_3 x^2 + a_4) \dot{x} + x = 0.$$
(14)

We find that :

$$-Aw^{2}\cos wt - (a_{0}A^{5}\cos wt^{5} + a_{1}A^{4}\cos wt^{4} + a_{2}A^{3}\cos wt^{3} + a_{3}A^{2}\cos wt^{2} + a_{4}) (Aw\sin wt + A\cos wt) = 0,$$
$$A(1 - w^{2})\cos wt = a_{0}A^{6}w\cos wt^{5}\sin wt + a_{1}A^{2}w\cos wt^{4}\sin wt$$

$$+ a_2 A^4 w \cos w t^3 \sin w t + a_3 A^3 w \cos w t^3 \sin w t + a_4 A w \sin w t$$

$$= Aw \left(a_0 A^5 \cos wt^5 + a_1 A^4 \cos wt^4 + a_2 A^3 \cos wt^3 + a_3 A^2 \cos wt^2 + a_4 \right) \sin wt^4 + a_2 A^3 \cos wt^3 + a_3 A^2 \cos wt^2 + a_4 \right)$$

After the simplification, we find the second member equal at :

$$= a_0 \frac{A^6 w}{32} \sin 6wt + a_1 \frac{A^5 w}{16} \sin 5wt + (a_0 \frac{A^6 w}{8}) + a_2 \frac{A^4 w}{8} \sin 4wt + (3a_1 \frac{A^5 w}{16}) + a_3 \frac{A^5 w}{4} \sin 3wt + (a_0 \frac{A^6 w}{32} + a_2 \frac{A^4 w}{4}) \sin 2wt + (a_1 \frac{A^5 w}{8} + a_3 \frac{A^3 w}{4} + a_4 Aw) \sin wt$$

This equation can be satisfied for all t only if the coefficients of the different sinusoidal terms disappear. The term cos(wt) disappears if w = 1 and the coefficient of sin wt is zero as if we take

$$Aw\left(a_{1}\frac{A^{4}}{8} + a_{3}\frac{A^{2}}{4} + a_{4}\right) = 0 \qquad Aw \neq 0 \quad in \ order \ not \ to \ fall \ in \ trivial \ case.$$
$$a_{1}A^{4} + 2a_{3}A^{2} + 8a_{4} = 0. \ We \ put \ y = A^{2}, so \ \delta = 4 \ a_{3}^{2} - 32 \ a_{1}a_{4}.$$

If $\delta = 0 \rightarrow a_3^2 = 8a_1a_4$, then $y = \frac{-a_3}{a_1}$ so the system accepts a limited cycle of radius $A = \sqrt{\frac{-a_3}{a_1}}$ as $a_1, a_3 < 0, a_1 \neq 0$.

 $y_{1,2} = \frac{-4a_3 \pm \sqrt{\delta}}{2a_{11}}.$

If $\delta > 0 \rightarrow a_3^2 > 8a_1a_4$, we find two solutions :

So, the system accepts a limited cycle only if $A_{1,2} = \sqrt{y_{1,2}}$ provided that $y_{1,2} > 0$. If $\delta < 0$, it is not accepts any limited cycle.

5thstep

We choose $f_m(x) = a_0 x^5 + a_1 x^4 + a_2 x^3 + a_3 x^2 + a_4 x + a_5$, $g_n(x) = x$.

The system becomes :

$$\begin{cases} \dot{x} = y \\ \dot{y} = -(a_0 x^5 + a_1 x^4 + a_2 x^3 + a_3 x^2 + a_4 x + a_5)y - x \end{cases}$$

In first approximation, the general solution is

$$x = C \cos(t - \alpha).$$

Where C and α are arbitrary. We study how the presence of the term with ε affects on the solution

$$x = C cost.$$

By giving what can be a first reasonable approximation, we try

$$x = A \cos wt$$
..

Where ω is a constant close to 1.

$$\begin{cases} \dot{x} = -A \, w \sin wt \\ \ddot{x} = -Aw^2 \cos wt \end{cases}$$

We replace in the following differential equation

$$\ddot{x} + (a_0 x^5 + a_1 x^4 + a_2 x^3 + a_3 x^2 + a_4 x + a_5) \dot{x} + x = 0.$$
(15)

We find that :

$$-Aw^{2}\cos wt - (a_{0}A^{5}\cos wt^{5} + a_{1}A^{4}\cos wt^{4} + a_{2}A^{3}\cos wt^{3} + a_{3}A^{2}\cos wt^{2} + a_{4}) (Aw\sin wt + A\cos wt) = 0,$$

 $\begin{aligned} A(1 - w^2)\cos wt &= a_0 A^6 w \, \cos wt^5 \sin wt + a_1 A^5 w \, \cos wt^4 \sin wt \\ &+ a_2 A^4 w \cos wt^3 \, \sin wt + a_3 A^3 w \, \cos wt^2 \sin wt + a_4 A^2 w \, \sin wt + a_5 A w \, \sin wt, \end{aligned}$

$$= Aw (a_0 A^5 \cos wt^5 + a_1 A^4 \cos wt^4 + a_2 A^3 \cos wt^3 + a_3 A^3 \cos wt^2 + a_3 A^2 \cos wt^2 + a_4 A \cos wt + a_5) \sin wt.$$

After the simplification, we find the second member equal at :

$$a_{0}\frac{A^{6}w}{32}\sin 6wt + a_{1}\frac{A^{5}w}{16}\sin 5wt + (a_{0}\frac{A^{6}w}{8}) + a_{2}\frac{A^{4}w}{8})\sin 4wt + (3a_{1}\frac{A^{5}w}{16}) + (3a_{1}\frac{A^{5}w}{16}) + (3a_{1}\frac{A^{5}w}{16}) + (a_{2}\frac{A^{4}w}{4})\sin 3wt + (a_{0}\frac{A^{6}w}{32}) + (a_{1}\frac{A^{5}w}{8} + a_{3}\frac{A^{3}w}{4} + a_{5}Aw)\sin wt.$$

This equation can be satisfied for all t only if the coefficients of the different sinusoidal terms disappear. The term cos(wt) disappears if w = 1 and the coefficient of sin wt is zero as if we take

$$Aw\left(a_{1}\frac{A^{4}}{8} + a_{3}\frac{A^{2}}{4} + a_{5}\right) = 0 \qquad Aw \neq 0 \quad in \ order \ not \ to \ fall \ in \ trivial \ case.$$
$$a_{1}A^{4} + 2a_{3}A^{2} + 8a_{5} = 0. \ We \ put \ y = A^{2}, so \ \delta = 4 \ a_{3}^{2} - 32 \ a_{1}a_{5}.$$

If $\delta = 0 \rightarrow a_3^2 = 8a_1a_5$, then $y = \frac{-a_3}{a_1}$ so the system accepts a limited cycle of radius $A = \sqrt{\frac{-a_3}{a_1}}$ as $a_1, a_3 < 0, a_1 \neq 0$.

If $\delta > 0 \rightarrow a_3^2 > 8a_1a_5$, we find two solutions : $y_{1,2} = \frac{-4a_3 \pm \sqrt{\delta}}{2a_{11}}$.

So, the system accepts a limited cycle only if $A_{1,2} = \sqrt{y_{1,2}}$ provided that $y_{1,2} > 0$. If $\delta < 0$, it accepts any limited cycle.

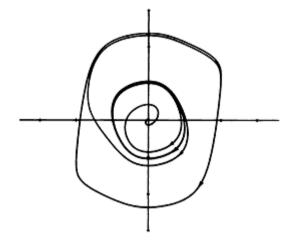


Figure : The two limit cycle of the Lienard system

8.Conclusion

In this paper, we proved the system can have at most 2 limit cycles. If f(x) is an odd polynomial of degree 5 then the probabilities that the Liénard equation for f(x) has at least 2 periodic solutions is greater than 47,23 % and that it has no periodic solution is greater than 34.54%.

References

- A.A. Andronov, Les cycles limites de Poincaré et la théorie des oscillations auto-entretenues, C. R. Acad. Sci. Paris 89 (1929) 559–561.
- [2]. C. Li, "Abelian integrals and limit cycles,"Qualitative Theory of Dynamical Systems, vol.11, no.1, pp. 111–128, 2012. [3] Briot, Ch. and J.Cl. Bouquet, Recherches sur les propriétés des fonctions définie par des équations différentielles, J. Ecole Imp. Polytech, 21 (1856).
- [3]. C. Christopher, "Estimating limit cycle bifurcations from center," in Differential Equations with Symbolic Computation, Trends in Mathematics, pp. 23–35, Birkhauser, Basel, Switzer-" land, 2005.
- [4]. Dulac, H, Solution d'un système d'équation diérentialles dans voisinage 133-197 de valeurs singuliéres.Bull. Soc. Math. France, 40(1912), 324-383.
- [5]. F. Dumortier, J. Llibre, J.C. Artés, Qualitative Theory of Planar Differential Systems, Universitext, Springer-Verlag, New York, 2006
- [6]. F. FENENICHE, M.S REZAOUI, Integrability and a Limit Cycle Solver for a Generalization of Polynomial Liénard Differential Systems, Journal of Applied Computer Science & Mathematics, Suceava, Vol.12, No. 26, 41-47, (2018).
- [7]. J. Llibre, D. Schlomiuk; The geometry of quadratic differential systems with a weak focus of third order, Canad.J. Math. 56 (2004), No. 2, 310-343.
- [8]. J. Llibre, R. Ramirez, M. Sadovskaia, On the 16th Hilbert problem for limit cycles on non-singular algebraic curves, J. Differential Equations 250 (2010) 983-999.
- [9]. Llibre, J., Zhang, X.: Limit cycles of the classical Liénard différential systems: a survey on the Lins Neto, de Melo and Pugh's conjecture, Expositiones Mathematicae Volume 35, Issue 3, 286-299 (2017)
- [10]. J. Li, H. S. Y. Chan, and K. W. Chung, "Investigations of bifurcations of limit cycles inZ2-equivariant planar vector fields of degree 5," International Journal of Bifurcation and Chaos in Applied Sciences and Engineering, vol. 12, no. 10, pp. 2137–2157, 2002.
- [11]. J.P.Demailly, Analyse numérique et équation différentiélles ,(1991).
- [12]. J. Li, H. S. Y. Chan, and K. W. Chung, "Bifurcations of limit cycles in a Z6-equivariant planar vector field of degree 5," Science in China. Series A, vol. 45, no. 7, pp. 817–826, 2002.
- [13]. Parker, T. S. and L. O. Chua, Practical Numerical Algorithms for Chaotic Systems.Springer-Verlag, New York.
- [14]. J. Sokulski, On the number of invariant lines of polynomial vector fields, Nonlinearity 9 (1996) 479–485.
- [15]. Regilene D. S. Oliveira, Alex C. Rezende, Dana Schlomiuk, Nicolae Vulpe; Geometric and algebraic classification of quadratic differential systems with invariant hyperbolas Electronic Journal of Differential Equations, Vol. 2017 (2017), No. 295, pp. 1122.ISSN: 1072-6691.)
- [16]. Sansone, G and R.Conti, Non-linear differential equation (translated by A.H.Diamond), Perg-amon Press, Macmillan, New York, 1964.
- [17]. S.CHARLES, Biologie Mathmatique et Modlisation ,Chapitre 3 : Fonctions de Lyapunov Notion de cycle limite page 31 BMM1 (scharles@biomserv.univ-lyon1.fr)01/02/2006.
- [18]. Sparrow C, Swinnerton-Dyer HPF. The Falkner–Skan equation. II. Dynamics and the bifurcations of P- and Q-orbits. J Differ Equations 2002;183(1):1–55.
- [19]. W. Xu and M. Han, "On the number of limit cycles of a Z4- equivariant quintic polynomial system," Applied Mathematics and Computation, vol. 216, no. 10, pp. 3022–3034, 2010.
- [20]. W. H. Yao and P. Yu, "Bifurcation of small limit cycles in Z5-equivariant planar vector fields of order 5," Journal of Mathematical Analysis and Applications, vol. 328, no. 1, pp. 400–413, 2007.
- [21]. X. Zhang, The 16th Hilbert problem on algebraic limit cycles, J. Differential Equations 251 (2011) 1778-1789.
- [22]. Yan-Qian Ye, Theory of limit cycles., American Mathematical Society 1986. Volume 66., Canad.J. Math. 56 (2004), No. 2, 310343.
- [23]. Y. Wu, L. Tian, and Y. Hu, "On the limit cycles of a Hamiltonian underZ4-equivariant quintic perturbation," Chaos, Solitons and Fractals, vol. 33, no. 1, pp. 298–307, 2007.
- [24].Zhang Z. Proof of the uniqueness of the limit cycles of generalized Lienard equation Appl. 1986. Anal. 23.63-76.